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RESEARCH ARTICLE

Mathematical Model of the Process of Ultrasonic wave Propagation in a Relax Environment with its Given Profiles at three Time Moments

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Abstract:

Objective:

The process of ultrasound oscillations in a relaxed environment, provided that the profiles of the acoustic wave at three time moments are known, is modeled by a three-point problem for the partial differential equation of the third order in time. This equation as a partial case contains a hyperbolic equation of the third order, which is widely used in ultrasound diagnostics.

Methods:

The differential-symbol method is applied to study a three-point in-time problem. The advantage of this method is the possibility to obtain a solution of the problem only through operations of differentiation.

Results:

We propose the formula to construct the analytic solution of the problem, which describes the process of ultrasound oscillations propagation in a relax environment. Due to this, the profile of the ultrasonic wave is known at any time and at an arbitrary point of space. The class of quasipolynomials is distinguished as a class of uniqueness solvability of a three-point problem.

Conclusion:

Using the proposed method, it is possible to analyze the influence of the main parameters of ultrasound diagnostics problems on the propagation of acoustic oscillations in a relaxed environment. The research example of a specific three-point problem is given.

Keywords: Mathematical model, Ultrasonic oscillations, Three-point in time problem, Differential-symbol method, Ultrasound diagnostics, Partial differential equation.

	Article History	Received: December 13, 2020	Revised: July 2, 2021	Accepted: August 21, 2021
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1. INTRODUCTION

To describe the processes of various natures, there are many models which are considered basic parameters of the process and they can effectively be investigated by mathematical methods.

Moreover, the same type of models is often used in completely different fields of knowledge. For example, in modeling biomechanical and medical processes [1 - 3], developed models for problems of hydromechanics and gas dynamics are used [4, 5].

One of the important areas of mathematical modeling is the simulation of processes by differential equations, which is used, in particular, in modeling the processes of wave propagation. Thus it is possible to predict their behavior depending on propagation conditions. In particular, the property of ultrasonic waves to change the speed of their propagation and absorption with any changes in the environment is taken into account. This property, as well as the reflection of the ultrasonic wave at the boundaries of different environments in the human body, are the basis of the ultrasound diagnostics, which is one of the most informative methods of non-invasive diagnosis in medicine. It is widely used to diagnose the work of organs and obtain their three-dimensional images, accelerate metabolic processes in the body and destroy various tumors.

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Today, more and more mathematical models of mechanical, biomedical and geophysical processes are used not only in partial differential equations of the second order in time, but also the third and fourth orders in time [6 - 9]. For studying these models, there are applied numerical, qualitative and asymptotic methods [10, 11].

The Cauchy problem for hyperbolic equation of a thirdorder in time of the form

$$\left[\tau \frac{\partial^3}{\partial t^3} - \tau c_1^2 \Delta \frac{\partial}{\partial t} + \frac{\partial^2}{\partial t^2} - c_2^2 \Delta \right] u(t, y) = 0, \quad y = (y_1, y_2, y_3) \in \mathbf{R}^3,$$
(1)

is investigated [12 - 14]. The problems of ultrasonic diagnostics contain the Eq. (1) in which u(t, y) is dynamic pressure, Δ is three-dimensional Laplace operator, constants c_1 and c_2 are limiting phase velocities of sound, τ is relaxation time.

In addition to the Cauchy problem for partial differential equations with initial conditions, problems with other time conditions are also intensively investigated. In particular, there are multipoint problems in which values of the unknown solution are given at different times. Problems with multipoint time conditions in unbounded domains are studied [15, 16], and in bounded domains with additional conditions of spatial variables are also studied [17 - 20].

Multipoint in time problems has an important physical interpretation and these are problems in which the states of the process are given at n points of time. In contrast to the Cauchy problem, these problems require detailed study because the kernel of multipoint in time problems is usually non-trivial [21]. Therefore, incorrect problems with multipoint conditions in time for partial differential equations demand new research methods. The differential-symbol method is sufficiently effective among the methods of research of multipoint problems in unbounded domains [22 - 25].

2. METHODS

In the present paper, the differential-symbol method is applied. This method has been successfully used to solve the problems for partial differential equations with time-variable conditions in unbounded domains. Such problems are the Cauchy problem as well as the problem with local and nonlocal multipoint in time conditions. To solve these problems, the Fourier transform method, which allows us to find solutions in integral form, in particular, in the form of improper integrals of given functions, is traditionally used. In contrast to the Fourier transform, the differential-symbol method makes it possible to construct solutions of the previous problems using differentiation operations. If the right-hand sides of conditions (for example, the initial functions in the Cauchy problem) are quasipolynomials $\varphi_i(x)$ then the solutions of the problem are represented in a form of actions of differential expressions ∂_

 $\varphi_i\left(\frac{\partial v}{\partial v}\right)$ on certain functions of a parameter or a vectorparameter v. After the actions of these differential expressions, the parameter v is assumed to be zero.

3. RESULTS

3.1. Posing of the Problem

Eq. (1) of ultrasound oscillations by replacement $x = (c_1 \tau)^{-1} y$, $\alpha = (c_2)^{-1} c_1$ is transformed into the one-parameter (parameter α $\hat{1}$ (0, 1)) hyperbolic equation

$$\left[\frac{\partial^3}{\partial t^3} - \Delta \frac{\partial}{\partial t} + \frac{\partial^2}{\partial t^2} - \alpha \Delta\right] u(t, x) = 0, \quad (t, x) \in (0, +\infty) \times \mathbf{R}^3.$$
 (2)

The mathematical model of the process of ultrasound wave propagation which contains Eq. (2) and the profile of oscillation is given at three equidistant moments of time t = jh, where $j\hat{1} J = \{0,1,2\}, h > 0$ (Eq. 3):

$$u(jh, x) = f_j(x), \quad j \in J, \quad x \in \mathbf{R}^3,$$
(3)

is considered [13]. In the presented paper, we obtain more general results.

Let us consider the problem for the partial differential equation of the third order (Eq. 4).

$$\left[\frac{\partial^3}{\partial t^3} + a\left(\frac{\partial}{\partial x}\right)\frac{\partial^2}{\partial t^2} + b\left(\frac{\partial}{\partial x}\right)\frac{\partial}{\partial t} + c\left(\frac{\partial}{\partial x}\right)\right]u(t,x) = 0, \qquad (t,x) \in (0,\infty) \times \square^3,$$
(4)

Which describes not only ultrasound waves as Eq. (2), but also different waves depending on the operator coefficients $\begin{pmatrix} \partial \\ \end{pmatrix}$, $\begin{pmatrix} \partial \\ \end{pmatrix}$

$$a\left(\frac{\partial}{\partial x}\right), \ b\left(\frac{\partial}{\partial x}\right) \text{ and } \ c\left(\frac{\partial}{\partial x}\right) \text{ , where}$$
$$a\left(\frac{\partial}{\partial x}\right) = \sum_{|s|\leq 2} a_s \frac{\partial^s}{\partial x^s}, \ b\left(\frac{\partial}{\partial x}\right) = \sum_{|s|\leq 2} b_s \frac{\partial^s}{\partial x^s}, \ c\left(\frac{\partial}{\partial x}\right) = \sum_{|s|\leq 3} c_s \frac{\partial^s}{\partial x^s},$$
(5)

 $\frac{\partial^{s_1}}{\partial x^{s_1}} = \frac{\partial^{s_1}}{\partial x_1^{s_1}} \frac{\partial^{s_2}}{\partial x_2^{s_2}} \frac{\partial^{s_3}}{\partial x_1^{s_1}}, \ s = (s_1, s_2, s_3), \ , \ |s| = s_1 + s_2 + s_3, \ a_s, \ b_s, \ c_s \text{ are complex numbers (Eq. 5).}$

The profiles of wave are given at time moments t, t_1 , t_2 by conditions (Eq. 6):

$$u(t_j, x) = g_j(x), \qquad j \in J, \quad x \in \square^3.$$
 (6)

Let us find a solution to problem (4)-(6) using the differential-symbol method [22].

According to Eq. (4), the ordinary differential equation with parameter μ has the form (Eq. 7):

$$\left[\frac{d^{3}}{dt^{3}} + a(\mu)\frac{d^{2}}{dt^{2}} + b(\mu)\frac{d}{dt} + c(\mu)\right]T(t,\mu) = 0,$$
(7)

where

To simplify the form of formulas, we will write $a = a(\mu)$, $b = b(\mu)$, $c=c(\mu)$. Let us denote the roots of the characteristic (cubic) equation (Eq. 8):

$$\lambda^3 + a\lambda^2 + b\lambda + c = 0 \tag{8}$$

by $\lambda_1 = \lambda_1(\mu)$, $\lambda_2 = \lambda_2(\mu)$, $\lambda_3 = \lambda_3(\mu)$. These roots have the following image (Eq. 9):

$$\lambda_k = l_k - \frac{a}{3}, \quad l_k = -\frac{\xi^k \chi + (\xi^k \chi)^{-1} \chi_0}{3}, \quad k = 1, 2, 3,$$
 (9)

Where
$$\chi_0 = a^2 - 3b$$
, $\chi_1 = 2a^3 - 9ab + 27c$, $\chi = \left(\frac{\chi_1 - \sqrt{\chi_1^2 - 4\chi_0^3}}{2}\right)^{1/3}$

and the cubic root of the one $\xi = (-1 + i\sqrt{3})/2$ belongs to the second quarter.

If the determinant $D = 18abc - 4a^3c + a^2b^2 - 4b^3 - 27c^2$

Of the polynomial $\lambda^3 + a\lambda^2 + b\lambda + c$ does not equal zero, then Eq. (8) has simple roots, if D=0, then the roots (9) are multiples. in particular

$$\lambda_1 = \frac{4ab - 9c - a^3}{a^2 - 3b}, \quad \lambda_2 = \lambda_3 = \frac{9c - ab}{2(a^2 - 3b)}$$

for $a^2 \neq 3b$ and $\lambda_1 = \lambda_2 = \lambda_3 = -a/3$ for $a^2 3b$.

For real numbers a, b, c if $D \ge 0$ there are real roots (if D >0 they are different) and if D < 0 we get real and two complex roots.

The geometric interpretation of the real roots is given in Fig. (1), where

$$\operatorname{Re}(\sigma + r^{3}e^{i3\varphi}) = \{S + r\cos\varphi, S + r\cos(\varphi + 2\pi/3), S + r\cos(\varphi + 4\pi/3)\}$$

is the set of roots and $\sigma = S + i \operatorname{Im} \sigma \in \mathbb{C}$, $\varphi \in [0, 2\pi/3)$, $r \ge 0$ (we have a double root $S \mp r/2$ for $\varphi = 0$, $\varphi = \pi/3$ and a triple root *S* for r = 0).



Fig. 1. Image of real roots $\lambda_1, \lambda_2, \lambda_3$.

For the equation of the third order (7) with the coefficients, which is integered by the vector-parameter μ , we write the normal fundamental at the point t₀ system of the solutions $\{\overline{T}_0(t,m), \overline{T}_1(t,m), \overline{T}_2(t,m)\}$ which is defined uniquely. In addition, we find the fundamental system $\{T_0(t,\mu), T_1(t,\mu), T_2(t,\mu)\}$ of solutions of Eq. (7), which satisfy conditions

$$T_{k}(t_{j},\mu) = \begin{cases} 1, & k = j, \\ 0, & k \neq j, \end{cases} \quad k, j \in J, \quad \mu \in M.$$
 (10)

The system $\{T_0(t,\mu), T_1(t,\mu), T_2(t,\mu)\}$ can be constructed only for vectors $\mu \in M \subseteq C^3$ for which the following condition

$$d(m)^{\circ} \det\left(\overline{T}_{k}(t_{j},m)\right)_{k,j\hat{l}} J^{-1} 0$$
(11)

is realized.

Therefore,

$$M = \{ m\hat{I} \ \mathbf{C}^3 : \ d(m)^1 \ 0 \}.$$
 (12)

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Set M, which is defined by formula (12), is divided in three sets $M = M_0 \cup M_1 \cup M_2$, moreover for vectors $\mu \in M_0$ roots of Eq (8) are simple, for $\mu \in M_1$ there are simple and double roots among the roots of Eq. (8) and a triple root for vectors $\mu \in M_2$.

$$G = \begin{bmatrix} G_{11}^{i} & G_{1233} & G \\ G_{21}^{i} & G_{2233} & G \\ G_{31}^{i} & G_{3233} & G \end{bmatrix} , \text{ which has}$$

the
$$G = \begin{bmatrix} e^{l_2 l_1 + l_3 l_2} - e^{l_1 l_2 + l_3 l_1} & e^{l_1 l_2 + l_3 l_1} - e^{l_1 l_1 + l_3 l_2} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_2 + l_3 l_1} \\ e^{l_2 l_2 + l_3 l_0} - e^{l_2 l_1 + l_3 l_2} & e^{l_1 l_1 + l_3 l_2} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_2 + l_2 l_1} \\ e^{l_2 l_0 + l_3 l_1} - e^{l_2 l_1 + l_3 l_0} & e^{l_1 l_1 + l_3 l_2} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_1 + l_2 l_2} \\ e^{l_2 l_0 + l_3 l_1} - e^{l_2 l_1 + l_3 l_0} & e^{l_1 l_1 + l_3 l_0} - e^{l_1 l_0 + l_3 l_1} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_1 + l_2 l_2} \\ e^{l_2 l_0 + l_3 l_1} - e^{l_2 l_1 + l_3 l_0} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_0 + l_3 l_1} \\ e^{l_1 l_1 - l_2 l_2 l_1 + l_2 l_2} & e^{l_3 l_1} \\ G^* = \begin{pmatrix} (t_2 - t_1) e^{l_2 (l_1 + t_2)} & t_2 e^{l_1 l_1 + l_2 l_2} - t_1 e^{l_1 l_2 + l_2 l_1} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_0 + l_2 l_2} \\ (t_1 - t_2) e^{l_2 (l_0 + t_2)} & t_2 e^{l_1 l_1 + l_2 l_2} - t_0 e^{l_1 l_2 + l_2 l_0} & e^{l_1 l_2 + l_2 l_0} - e^{l_1 l_0 + l_2 l_2} \\ (t_1 - t_0) e^{l_2 (l_0 + t_1)} & t_0 e^{l_1 l_1 + l_2 l_2} - t_0 e^{l_1 l_2 + l_2 l_1} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_1 + l_2 l_2} \\ (t_1 - t_0) e^{l_2 (l_0 + t_1)} & t_0 e^{l_1 l_1 + l_2 l_2} - t_0 e^{l_1 l_2 + l_2 l_1} & e^{l_1 l_1 + l_2 l_2} - e^{l_1 l_1 + l_2 l_2} \\ \end{bmatrix}$$

If $\mu \in M_0$ (simple roots), then

$$T_{0}(t,\mu) = \frac{G_{11}e^{l_{1}t} + G_{12}e^{l_{2}t} + G_{13}e^{l_{3}t}}{G_{11}e^{l_{1}t_{0}} + G_{12}e^{l_{2}t_{0}} + G_{13}e^{l_{3}t_{0}}}e^{a(t_{0}-t)/3},$$

$$T_{1}(t,\mu) = \frac{G_{21}e^{l_{1}t} + G_{22}e^{l_{2}t} + G_{23}e^{l_{3}t}}{G_{21}e^{l_{1}t} + G_{22}e^{l_{2}t} + G_{23}e^{l_{3}t}}e^{a(t_{1}-t)/3},$$

$$T_{2}(t,\mu) = \frac{G_{31}e^{l_{1}t} + G_{32}e^{l_{2}t} + G_{33}e^{l_{3}t}}{G_{31}e^{l_{1}t} + G_{32}e^{l_{2}t} + G_{33}e^{l_{3}t}}e^{a(t_{2}-t)/3}.$$
(13)

For the case $\mu \in M_1$ (root λ_1 is simple and λ_2 is double), then

$$T_{0}(t,\mu) = \frac{G_{11}^{e_{11}}e^{l_{11}} + (G_{12}^{e_{12}} + G_{13}^{e_{13}}t)e^{l_{21}}}{G_{11}^{e_{11}}e^{l_{10}} + (G_{12}^{e_{22}} + G_{13}^{e_{13}}t_{0})e^{l_{210}}}e^{a(t_{0}-t)/3},$$

$$T_{1}(t,\mu) = \frac{G_{21}^{e_{11}}e^{l_{11}} + (G_{22}^{e_{22}} + G_{23}^{e_{32}}t)e^{l_{21}}}{G_{21}^{e_{11}}e^{l_{11}} + (G_{22}^{e_{22}} + G_{23}^{e_{33}}t)e^{l_{21}}}e^{a(t_{1}-t)/3},$$

$$T_{2}(t,\mu) = \frac{G_{31}^{e_{11}}e^{l_{11}} + (G_{32}^{e_{22}} + G_{33}^{e_{33}}t)e^{l_{21}}}{G_{31}^{e_{11}}e^{l_{12}} + (G_{32}^{e_{32}} + G_{33}^{e_{33}}t)e^{l_{21}}}e^{a(t_{2}-t)/3}.$$
(14)

Eqs (13, 14) are correct (determinants are nonzero) and because 1,2,3 we have т
$$\begin{split} \sum_{j=1}^{3} G_{mj} e^{l_{j}t_{m-1}} &= \delta(\mu) e^{(t_{0}+t_{1}+t_{2})a/3} \neq 0, \qquad \mu \in M_{0}, \\ G_{m1}^{*} e^{l_{i}t_{m-1}} + (G_{m1}^{*}+G_{m1}^{*}t_{m-1}) e^{l_{2}t_{m-1}} &= \delta(\mu) e^{(t_{0}+t_{1}+t_{2})a/3} \neq 0, \quad \mu \in M_{1}. \end{split}$$

For the case $\mu \in M_2$ (triple root) formulas of the elements of

the fundamental system of solutions of Eq. (7), which satisfy conditions (10) are simplified and take the form

$$T_{0}(t,\mu) = \frac{(t-t_{1})(t-t_{2})}{(t_{0}-t_{1})(t_{0}-t_{2})}e^{a(t_{0}-t)/3},$$

$$T_{1}(t,\mu) = \frac{(t-t_{0})(t-t_{2})}{(t_{1}-t_{0})(t_{1}-t_{2})}e^{a(t_{1}-t)/3},$$

$$T_{2}(t,\mu) = \frac{(t-t_{0})(t-t_{1})}{(t_{2}-t_{0})(t_{2}-t_{1})}e^{a(t_{2}-t)/3}.$$
(15)

We introduce the quasipolynomial functions of the form

$$g(x) = Q(x)e^{nx}, \quad n \times x = n_1x_1 + n_2x_2 + n_3x_3,$$
 (16)

For the set (12), K_M is a class of functions which consists of the quasipolynomials of variables x_1 , x_2 , x_3 which can be represented as a finite sum of quasipolynomials of the form (15) with the different vectors. Also, we suppose that the zero quasipolynomial g(x) = 0 belongs to K_M .

For each quasipolynomial of the form (15), the differential expression $s\left(\frac{\partial}{\partial \mu}\right)$ can be obtained by replacing the vector *x* with the vector-derivative $\frac{\partial}{\partial \mu}$. The differential expression $Q\left(\frac{\partial}{\partial \mu}\right)e^{n\frac{\partial}{\partial \mu}\cdot n\cdot \frac{\partial}{\partial \mu}}$ of the infinite order acts on the analytic function $\gamma(\mu)$ at the neighborhood of m=n by the formula

$$\mathcal{Q}\left(\frac{\partial}{\partial\mu}\right)e^{\nu_1\frac{\partial}{\partial\mu_1}+\nu_2\frac{\partial}{\partial\mu_2}+\nu_3\frac{\partial}{\partial\mu_3}}\gamma(\mu) \equiv \left\{\mathcal{Q}\left(\frac{\partial}{\partial\mu}\right)\gamma(\mu)\right\}\Big|_{\mu=0}$$

Similarly, we determine the effect of differential expressions of infinite order on the analytic functions for more general quasipolynomials of the class K_M . We also assume that

$$g\left(\frac{\partial}{\partial\mu}\right)\gamma(\mu) \equiv 0 \text{ for } g = 0.$$

3.2. Main result

For arbitrary functions $g_0(x)$, $g_1(x)$, $g_2(x)$ from the class K_M there is a unique solution of the problem (4), (6), which is represented by the formula

$$u(t,x) = \sum_{k=0}^{2} g_k \left(\frac{\partial}{\partial \mu} \right) \left\{ T_k(t,\mu) e^{\mu x} \right\} \bigg|_{\mu=0}, \qquad (17)$$

moreover, $u(t,.) \in K_M$ for all $t \in (0,\infty)$.

In Eq. (17) we assume O = (0,0,0)
$$\frac{\partial}{\partial \mu} = \left(\frac{\partial}{\partial \mu_1}, \frac{\partial}{\partial \mu_2}, \frac{\partial}{\partial \mu_3}\right)$$

The fact that function (17) satisfies the partial differential Eq. (4) follows from the commutativity of the differentiation operators $\frac{\partial}{\partial t}$, $\frac{\partial}{\partial x}$, $\frac{\partial}{\partial \mu}$ and that the functions $T_0(t,m)$, $T_1(t,m)$, $T_2(t,m)$ satisfy the ordinary differential Eq. (7).

For the function of the form (17) three-point conditions (6) are carried out, since $T_0(t,m), T_1(t,m), T_2(t,m)$ satisfy conditions (10) and the equalities

$$g_k\left(\frac{\partial}{\partial\mu}\right)e^{\mu x}\Big|_{\mu=0} = g_k(x), \ k \in J$$

are realized.

The solution of three-point problem (4), (6) in the class of quasipolynomials for each fixed t ϵ (0, ∞) belongs to K_M, is unique. This is proved by differential-symbol method (see, for example [23],). Choice of the functions $g_0(x)$, $g_1(x)$, $g_2(x)$ from class K_M is significant.

Obviously, if the functions $g_0(x)$, $g_1(x)$, $g_2(x)$ in conditions (6) belong to the set K_M , then according to Eq. (17), the solutions of problem (4), (6) are quasipolynomials.

Numerous studies [26 - 30] have been devoted to the construction of solutions of some partial differential equations and boundary value problems for these equations in the form of quasipolynomials.

Thus, the process of oscillation propagation, in particular, ultrasound wave in a relaxed environment with given wave

The process is modeled by the problem (4), (6) for the partial differential Eq. (4) of the third order in time with threepoint time conditions (6). The method of constructing the solution of problem (4), (6) is indicated. The class of quasipolynomials KM is distinguished as the class of existence of the unique solution of problem (4), (6). Eq. (4) belongs to partial differential equations, which are widely used in problems of ultrasound diagnostics.

profiles at three points at time is investigated.

3.3. Example of Research of the Process of Ultrasound wave Propagation with its given Profiles at three Moments of Time

Let us investigate the process of acoustic oscillations using the proposed method for a specific example.

Example. We study the process of propagation of an ultrasound wave in relax environment, which is described by problem (4), (6) with the following parameters:

$$a\left(\frac{\partial}{\partial x}\right) = 1, \ b\left(\frac{\partial}{\partial x}\right) = -\Delta, \ c\left(\frac{\partial}{\partial x}\right) = -\frac{1}{9}\Delta,$$

$$g_1(x) = \sin\frac{x_1 + x_2 - x_3}{3}, \ g_0(x) = 0, \ g_2(x) = 0,$$

$$t_0 = 0, \ t_1 = 1, \ t_2 = 2.$$

The profile of wave at the moment of time t=1 is a sum of two quasipolynomials of the form (16) that is

$$g_1(x) = \sin \frac{x_1 + x_2 - x_3}{3} = \frac{1}{2}e^{\frac{x_1 + x_2 - x_3}{3}i} - \frac{1}{2}e^{-\frac{x_1 + x_2 - x_3}{3}i},$$

in which the polynomials are $\frac{1}{2}$ and $-\frac{1}{2}$ respectively, then the vector-parameters in the exponent are equal to $\nu = \left(\frac{i}{3}, \frac{i}{3}, -\frac{i}{3}\right)$

and -v. These vectors v and -v belong to the set M_2 and the corresponding root of Eq. (8) is triple:

$$\lambda_1 = \lambda_2 = \lambda_3 = -\frac{a}{3} = -\frac{1}{3}.$$

Function (14) take the form

$$\begin{split} T_0(t,\pm\nu) &= \frac{(t-1)(t-2)}{2}e^{-\frac{t}{3}},\\ T_1(t,\pm\nu) &= -t(t-2)e^{\frac{1-t}{3}},\\ T_2(t,\pm\nu) &= \frac{t(t-1)}{2}e^{\frac{2-t}{3}}. \end{split}$$

Since $d(\pm n) = 2e^{-1} 0$, then $\pm n \hat{1} M$

The solution of problem (4), (6) with given parameters can be found by Eq. (16):

$$u(t,x) = \sin\left[\frac{1}{3}\left(\partial_{\mu_{1}} + \partial_{\mu_{2}} - \partial_{\mu_{3}}\right)\right]\left\{T_{1}(t,\mu)e^{\mu \cdot x}\right\}\Big|_{\mu=0}$$
$$= \frac{1}{2}\left\{T_{1}(t,\mu)e^{\mu \cdot x}\right\}\Big|_{\mu=\frac{i}{3}(1,1,-1)} - \frac{1}{2}\left\{T_{1}(t,\mu)e^{\mu \cdot x}\right\}\Big|_{\mu=-\frac{i}{3}(1,1,-1)}$$

We calculate



Fig. 2. Graphical dependence of the solution $u(t, \eta)$ on $t \hat{1} [0,20]$ and $\eta \in [0,15]$.

$$\left\{T_{1}(t,\mu)\right\}\Big|_{\mu=\frac{i}{3}(1,1,-1)}=\left\{T_{1}(t,\mu)\right\}\Big|_{\mu=-\frac{i}{3}(1,1,-1)}=t\left(2-t\right)e^{\frac{i}{3}(1-t)}.$$

Finally, we obtain the solution to the problem

$$u(t,x) = t \left(2-t\right) e^{\frac{1}{3}(1-t)} \sin \frac{x_1 + x_2 - x_3}{3}.$$

Let us investigate the solution of the problem on parallel planes $x_1 + x_2 - x_3 - 3h = 0$ of the space R³, where $h\hat{1}$ **R**. Then the solution of problem has the form (Eq. 18):

$$u(t,\eta) = \left(2t - t^2\right) e^{\frac{1}{3}(1-t)} \sin \eta$$
 (18)

and it is 2*p*-periodic function by variables *h*.

The graph of the function (18) of two variables t and h is given in Fig. (2).

Therefore, function (17) describes the periodic oscillations of the ultrasound wave with the period $T = 2\pi$ by the variable η . For $h = \frac{p}{6}$ and $h = \frac{p}{2}$ the amplitudes of oscillation are given in the following form:

$$A_{1}(t) = \left| \left(2-t\right)t \right| e^{\frac{1}{3}(1-t)}, \quad A_{2}(t) = \frac{1}{2} \left| \left(2-t\right)t \right| e^{\frac{1}{3}(1-t)}.$$

These amplitudes are shown in Fig. (3) by a solid and dashed line.

As mentioned above, the found solution of problem is unique in the class of quasipolynomials and belongs to K_M each fixed *t*. Therefore, the ultrasound wave oscillates in a limited range at any time and exponentially goes to zero for $t \rightarrow +\infty$ on planes which is parallel to the plane $x_1 + x_2 - x_3 =$ 0 and acquire zero values on the planes $x_1 + x_2 - x_3 = 2pk, k \hat{l} \mathbf{Z}$

4. DISCUSSION

It should be noted that in recent years the approaches in which differentiation operations are used instead of integration operations, which are more complex (see, for example, [31, 32]). In these works, instead of the integration operation, other formulas which provide for the operation of differentiation by parameter ε are proposed. In particular, the following formulas are given for calculating integrals of the first kind:



Fig. (3). Graphs of amplitudes of the oscillating process for $A_1(t) \ge A_2(t)$.

$$\int_{-\infty}^{0} f(x) dx = \lim_{\varepsilon \to 0} f\left(\frac{\partial}{\partial \varepsilon}\right) \frac{1}{\varepsilon},$$

$$\int_{-\infty}^{\infty} f(x) dx = \lim_{\varepsilon \to 0} \left[f\left(\frac{\partial}{\partial \varepsilon}\right) + f\left(-\frac{\partial}{\partial \varepsilon}\right) \right] \frac{1}{\varepsilon}.$$

To calculate the Laplace transform $L[f](x) = \int e^{-xy} f(y) dy$ of the function f(x) there is proposed the following formula

$$L[f](x) = f\left(-\frac{\partial}{\partial x}\right)\frac{1}{x}$$

in which a differentiation by the variable x is used.

Therefore, the method which is proposed in this article is a new and effective method of solving the problem with threepoint in time conditions.

CONCLUSION

The process of ultrasound wave propagation in the environment with relaxation under the condition of setting the wave profile at three points of time is simulated by the threepoint problem for a partial differential equation of the third order in time.

A class of functions in which there is a unique solution to the problem has been established. It is a class of quasipolynomials with certain conditions on the index in exponents. An effective method for constructing a solution to the problem in this class is proposed. The process of acoustic oscillations is studied for a specific example. In the future, it is interesting to study acoustic oscillations under the influence of external forces, which leads to the study of the problem with three point in time conditions for a nonhomogeneous partial differential equation of the third order in time.

The proposed research results can be used in biomechanics and medicine, in particular, in the theory of ultrasound diagnostics.

ETHICS	APPROVAL	AND	CONSENT	ТО
PARTICIP	ATE			

Not applicable.

HUMAN AND ANIMAL RIGHTS

No humans and animals were used for studies that are the basis of this research.

CONSENT FOR PUBLICATION

Not applicable.

AVAILABILITY OF DATA AND MATERIALS

Not applicable.

FUNDING

None

CONFLICT OF INTEREST

The authors declare no conflict of interest, financial or otherwise.

ACKNOWLEDGEMENTS

Declared none.

REFERENCES

- [1] Kapur NJ. Mathematical models in biology and medicine. New Delhi: Affiliated East-West Press 1985.
- Horn MA, Simonett G, Webb GF. Mathematical models in medical [2] and health science. Vanderbilt University Press 1998.
- [3] Begun PI, Afonin PN. Modeling in biomechanics Moscow: Vysshaya shkola 2004
- Samarskii AA. The theory of difference schemes. CRC Press 2001. [4] [http://dx.doi.org/10.1201/9780203908518]
- Samarskii AA, Mikhailov AP. Principles of mathematical modelling: [5] Ideas, methods, examples. CRC Press 2013.
- Morrison JA. Wave propagation in rods of Voigt material and visco-[6] elastic materials with three-parameter models. Q Appl Math 1956; 14: 153-69
 - [http://dx.doi.org/10.1090/qam/78848]
- Rudenko OV, Soluyan SI. Theoretical Foundations of Nonlinear [7] Acoustics. Moscow: Nauka 1975.
- Lobanov AI, Starozhilova TK, Guriya GT. Numerical investigation of [8] pattern formation in blood coagulation. Math Model 1997; 9(8): 83-95.
- Agapov PI, Vasyukov AV, Petrov IB. Computer simulation of wave [9] processes in the integument of the brain during traumatic brain injury Processes and methods of information processing M. MFTI 2006; pp. 154-63
- Clarke JF, McChesney M. Dynamics of Relaxing Gases. London: [10] Butterworth's 1976.
- [11] Lick W. Wave propagation in real gases. Adv Appl Mech 1967; 10: 1-72.
- Renno P. The fundamental solution of a hyperbolic operator in three-[12] dimensional thermochemistry. Rend Accad Naz Sci XI Mem Mat 1979/1980; 4: 43-62.
- [13] Nytrebych Z, Il'kiv V, Malanchuk O. On the modeling process of ultrasonic wave propagation in a relaxation medium by the three-point in time problem. CEUR Workshop Proceedings (3rd International

Workshop on Informatics & Data-Driven Medicine (IDDM-2020). 2753: 72-81.

- [14] Varlamov V. Time estimates for the Cauchy problem for a third-order hyperbolic equation. Int J Math Math Sci 2003; 17: 1073-81. [http://dx.doi.org/10.1155/S0161171203204361]
- [15] Borok VM. Perel'man MA. Unique solution classes for a multipoint boundary value problem in an infinite layer. Izv Vys Uchebn Zaved Matematika 1973; 135(8): 29-34.
- Antypko II, Borok VM. Criterion of unconditional correctness of a [16] boundary value problem in a layer. Theory of Functions. Funct Anal Appl 1976; 26: 3-9.
- Incorrect boundary value problem for partial differential equations. [17] Kviv: Nauk. Dumka 1984.
- [18] Ilkiv V, Ptashnyk B. Problems for partial differential equations for nonlocal conditions. Metric approach to the problem of small denominators. Ukr Math J 2006; 58(12): 1847-75. [http://dx.doi.org/10.1007/s11253-006-0172-8]
- [19] Ptashnyk BI, Symotyuk MM. Multipoint problem for nonisotropic partial differential equations with constant coefficients. Ukr Math J 2003; 55(2): 293-310.

[http://dx.doi.org/10.1023/A:1025468413500]

- Kalenyuk PI, Volyans'ka II, Il'kiv VS, Nytrebych ZM. On the unique [20] solvability of a three-point problem for partial differential equation in a two-dimensional domain. J Math Sci 2020; 246(2): 170-87. [http://dx.doi.org/10.1007/s10958-020-04728-x]
- Malanchuk O, Nytrebych Z. Homogeneous two-point problem for [21] PDE of the second order in time variable and infinite order in spatial variables. Open Math 2017; 15(1): 101-10. [http://dx.doi.org/10.1515/math-2017-0009]
- [22] Nytrebych ZM. An operator method of solving the Cauchy problem for a homogeneous system of partial differential equations. J Math Sci 1996; 81(6): 3034-8. [http://dx.doi.org/10.1007/BF02362589]
- [23] Nytrebych ZM, Malanchuk OM. The differential-symbol method of solving the problem two-point in time for a nonhomogeneous partial differential equation. J Math Sci 2017; 227(1): 68-80. [http://dx.doi.org/10.1007/s10958-017-3574-2]
- [24] Nytrebych Z, Malanchuk O. The conditions of existence of a solution of the two-point in time problem for nonhomogeneous PDE. Int J Pure Appl Math 2019; 41: 242-50.
- [25] Nytrebych Z, Malanchuk O. The conditions of existence of a solution of the degenerate two-point in time problem for PDE. Asian-European Journ Math 2019; 12(3)1950037

[http://dx.doi.org/10.1142/S1793557119500372]

[26] Karachik V. Construction of polynomial solutions to some boundary value problems for Poisson's equation. Comput Math Math Phys 2011; 51(9): 1567-87.

[http://dx.doi.org/10.1134/S0965542511090120]

- Algazin OD. Exact solutions to the boundary-value problems for the [27] Helmholtz equation in a layer with polynomials in the right-hand sides of the equation and of the boundary conditions. Bulletin of the Moscow state regional University. Series. Physics-Mathematics 2020; 1:6-27.
- [28] Pedersen P. A basis for polynomial solutions for systems of linear constant coefficient PDE's. Adv Math 1996; 117: 157-63. [http://dx.doi.org/10.1006/aima.1996.0005]
- [29] Nytrebych ZM, Malanchuk OM. The differential-symbol method of constructing the quasi-polynomial solutions of two-point problem. Demonstr Math 2019; 52(1): 88-96. [http://dx.doi.org/10.1515/dema-2019-0010]

Hayman WK, Shanidze ZG. Polynomial solutions of partial

- [30] differential equations. Methods Appl Anal 1999; 6(1): 97-108. Kempf A, Jackson DM, Morales AH. How to (path-) integrate by [31]
- differentiating. J Phys Conf Ser 2015; 626(1)
- [32] Ding J, Tang E, Kempf A. Integration by differentiation: new proofs, methods and examples. J Phys A Math Theor 2017; 50(23) arxiv.org/abs/1610.09702v2

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